



Studies of Blazar emission using a spatially resolved SSC

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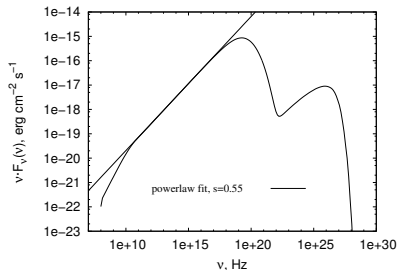
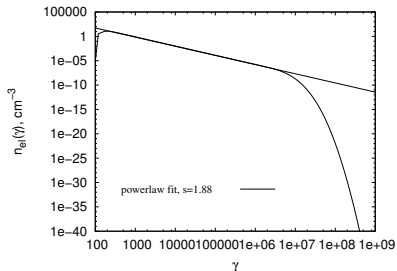
- ① A spatially resolved model
 - Why?
 - Acceleration at shock
 - Radiation
 - Photohadronics

- ② Results
 - SED of Mkn501 in 2009
 - Modelling of ultra short variability
 - Spectral index of the radio emission

- ③ Outlook
 - Hybrid acceleration
 - Oblique, relativistic shocks

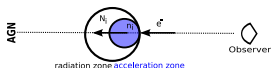


- get an electron distribution from somewhere (powerlaw, kinetic equation, shock simulations)
- assume these electrons to live in a small region within a jet (a so called blob)
- account for and compute any relevant radiative processes to gain the spectrum emitted from that region





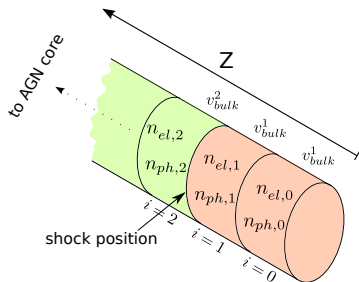
- modelling variability on timescales smaller than the light crossing time
- time dependent modelling of acceleration without artificial boundary between acceleration and radiation region
- in a hybrid model full treatment of intermediates (cooling, reacceleration(?)) is possible
- connection between the spectra of the central emission region ($\sim 10^{15}$ cm) and VLBI blobs ($\sim 10^{18}$ cm)
- \Rightarrow origin of VHE radiation from modelling correlation between VHE and radio regime
- one first step towards modelling of polarisation

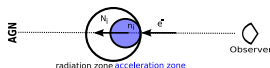


expand two-zone-model
to N-zone-model



move from nested
to aligned setup

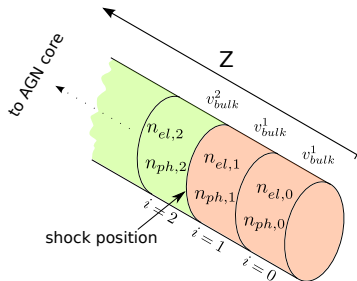




expand two-zone-model
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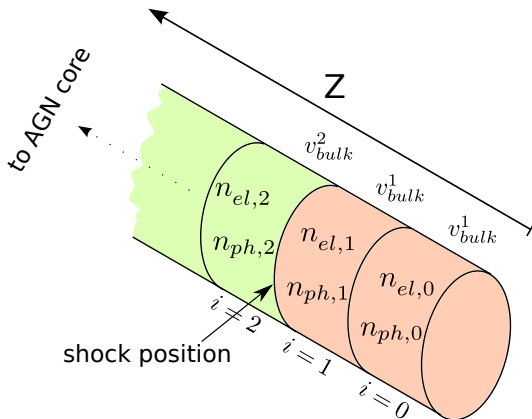
move from nested to
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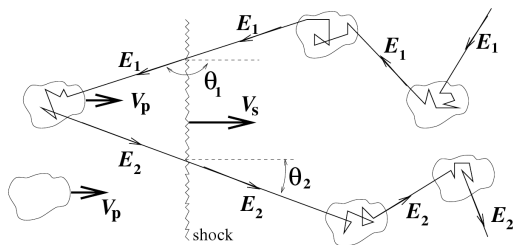


Not a multi
component model!



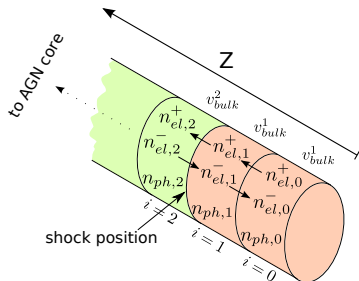
- divide simulation region into N slices in the direction parallel to shock normal
- each slice has local bulk speed, electron density, photon density (, magnetic field, radius)





mimic Fermi I acceleration

- ① advection between up- and downstream region
- ② pitch angle scattering that is isotropic in the local bulk frame



- split electron population into two half spheres; one moving downstream (n^+), the other moving upstream (n^-)
- advection in the shock frame:

$$\frac{\partial n}{\partial t} + \frac{v_{bulk} + \mu}{1 + v_{bulk}\mu} \frac{\partial n}{\partial x} = 0$$

- averaging over μ for each half-space
- scattering from one half space into the other (including boost)



time evolution of photon density

$$\frac{\partial N}{\partial t} = -c \cdot \kappa_{\nu,SSA} \cdot N + \frac{4\pi}{h\nu} \cdot (\epsilon_{\nu,IC} + \epsilon_{\nu,sync})$$

- $\kappa_{\nu,SSA}$ - Synchrotron Self Absorption coefficient
calculated using the Melrose Approximation
- $\epsilon_{\nu,IC}$ - changes due to invers compton scattering
full integration of photon and electron density using the Klein-Nishina cross section
- $\epsilon_{\nu,sync}$ - yields due to synchrotron radiation
integration of electron density using the Melrose approximation for a single electron spectrum

see Richter and Spanier 2012 for details



cooling (and maybe acceleration) of intermediates is relevant in AGNs

implementation

- using model by Hümmer et al. 2010
- full treatment of intermediate species
- e.g. calculation of realistic, flavour splitted neutrino spectra

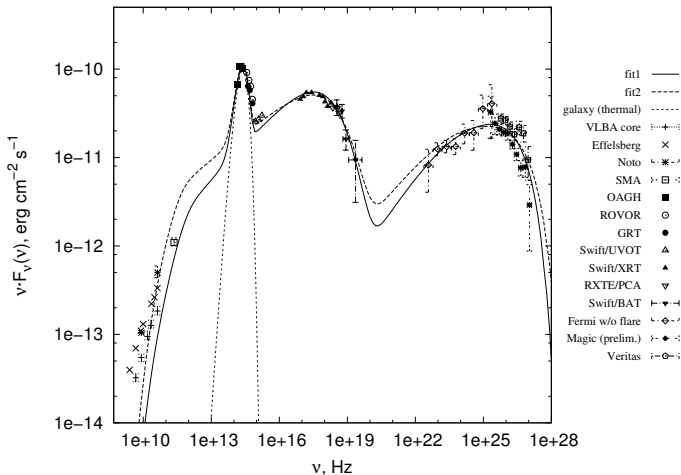
status

- arbitrary combinations of particle species and all relevant processes implemented
- problem: what's the correct treatment of electron and proton acceleration in a parallel shock? \Rightarrow see outlook

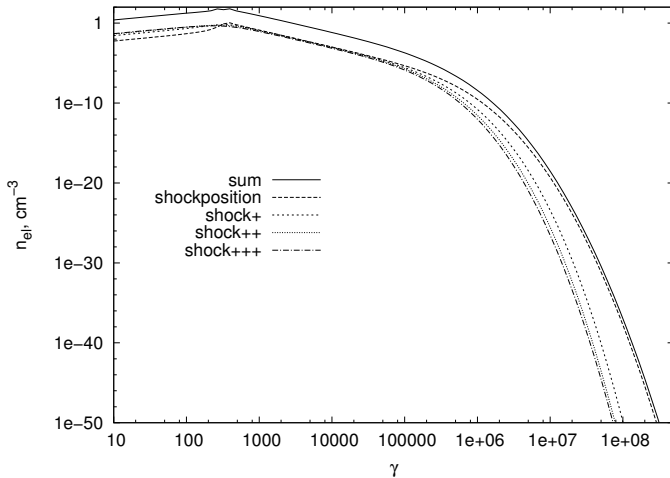


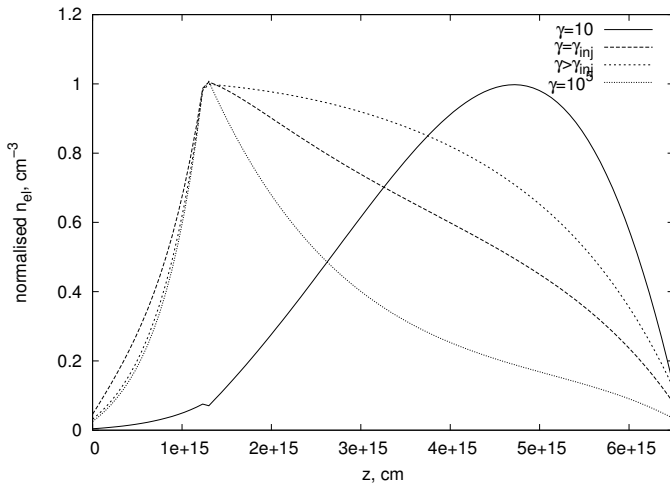
Results

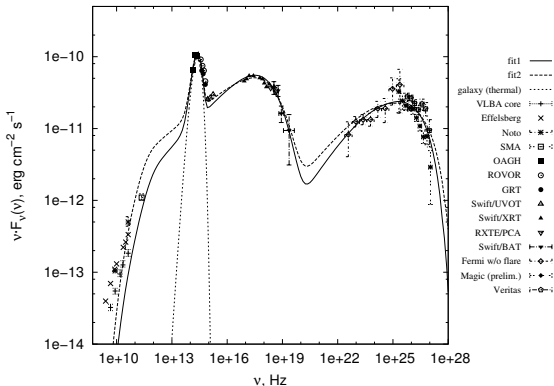
Steady state fit



Data taken from Abdo et al. (2011).







..yields

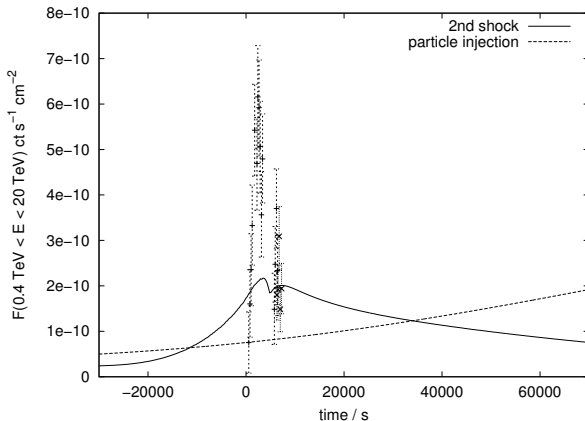
- fit1: $R = 6.5 \cdot 10^{15}$ cm and $\delta = 37$, hence $t_{lc} = 5856$ s
 - fit2: $R = 2.1 \cdot 10^{16}$ cm and $\delta = 47$, hence $t_{lc} = 14\,900$ s
- comparable parameters to one zone fits



- but variability is restricted by acceleration timescale

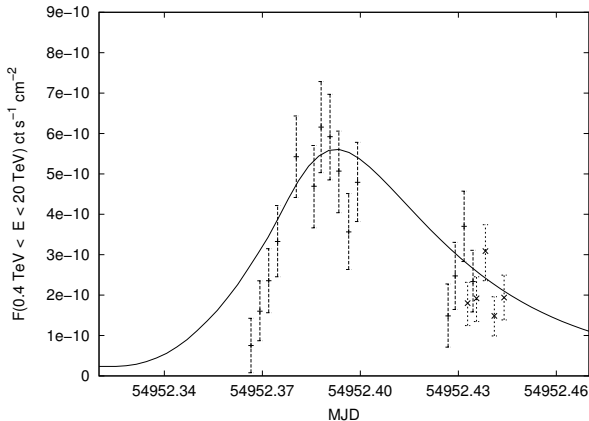
$$t_{var} \gtrsim \max(t_{lc}, t_{acc} \vee t_{cool})$$

- variability due to multiple shocks





- but variability is very similar to IC timescale



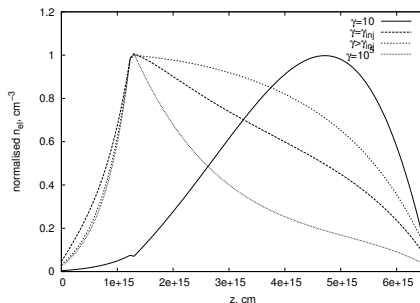
possible scenario summarised in Richter & Spanier (submitted)



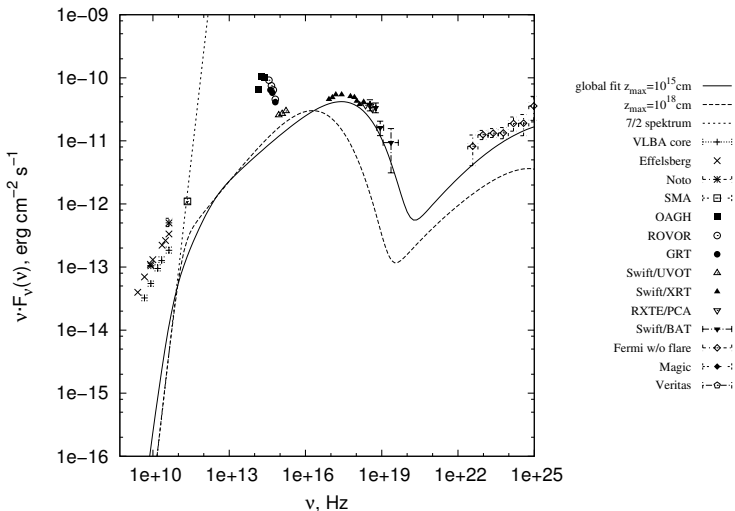
- naive spectral index: $F_\nu \propto \nu^{5/2}$ (self absorbed regime)
- in homogenous fits either ignored or circumvented by high γ_{inj} and γ_{min} , respectively (don't do the latter)



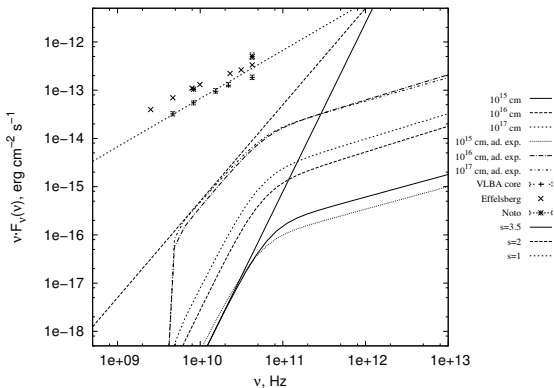
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⇒ far downstream there will be a powerlaw also for $\gamma < \gamma_{inj}$

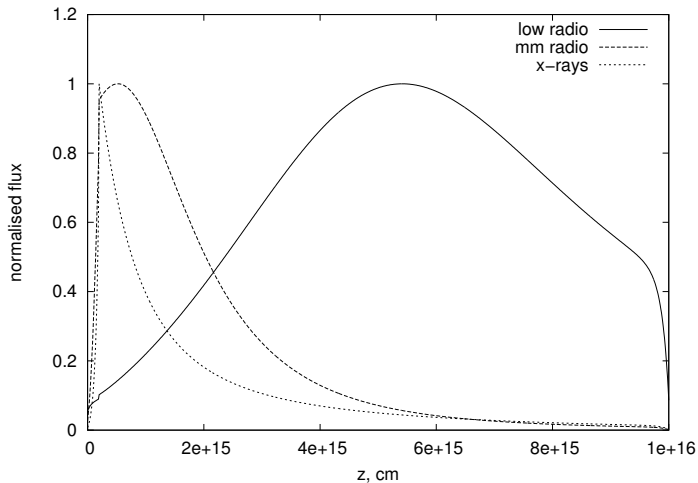


Effect of significant larger simulation region.



Effect of adiabatic expansion in the downstream.

- ⇒ yields opening angle of order $\tan(\alpha) \sim 1$
- ⇒ models in 2D or spherical geometry needed



Flux morphology.



Outlook



(parallel) shock escape time

- even homogenous models assume first order acceleration \Rightarrow shock is needed (as far as we know)
- since scattering assumed isotropic in plasma frame \Rightarrow escape time basically determined by shock speed and $D_{\mu\mu}$

consequences for acceleration timescale

- t_{esc} and t_{acc} have opposite dependence on $D_{\mu\mu}$, but we see powerlaws
- hence $D_{\mu\mu}$ (therefore acceleration efficiency) should be independent from energy and mass
- crucial for timescale predictions (see talk by M. Weidinger)
- identification of hadronic sources could give hints to the underlying turbulence
- oblique/relativistic jets? time dependent implementation?



Thank you

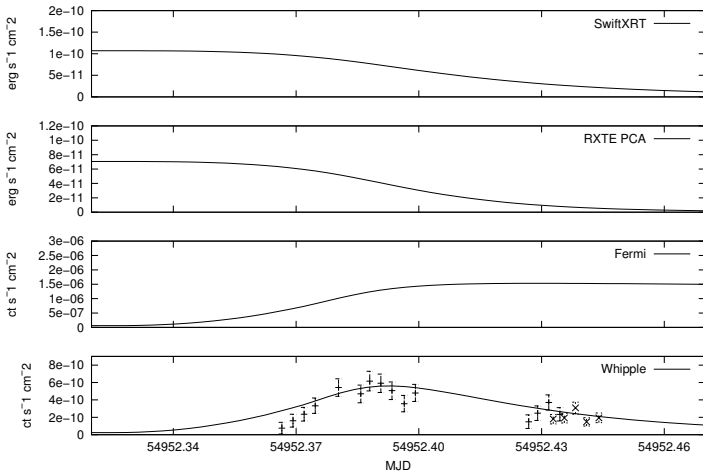


- A. A. Abdo, M. Ackermann, M. Ajello, A. Allafort, L. Baldini, J. Ballet, G. Barbiellini, M. G. Baring, D. Bastieri, K. Bechtol, and E. al. et al. Insights into the High-energy $\{\gamma\}$ -ray Emission of Markarian 501 from Extensive Multifrequency Observations in the Fermi Era. *The Astrophysical Journal*, 727(2):129, Feb. 2011. ISSN 0004-637X. doi: 10.1088/0004-637X/727/2/129. URL <http://stacks.iop.org/0004-637X/727/i=2/a=129?key=crossref.3481d95923711c6992ac11403c374cdf>.
- S. Hümmer, M. Rüter, F. Spanier, and W. Winter. Simplified Models for Photohadronic Interactions in Cosmic Accelerators. *The Astrophysical Journal*, 721(1232):630–652, Sept. 2010. doi: 10.1088/0004-637X/721/1/630. URL <http://iopscience.iop.org/0004-637X/721/1/630>.
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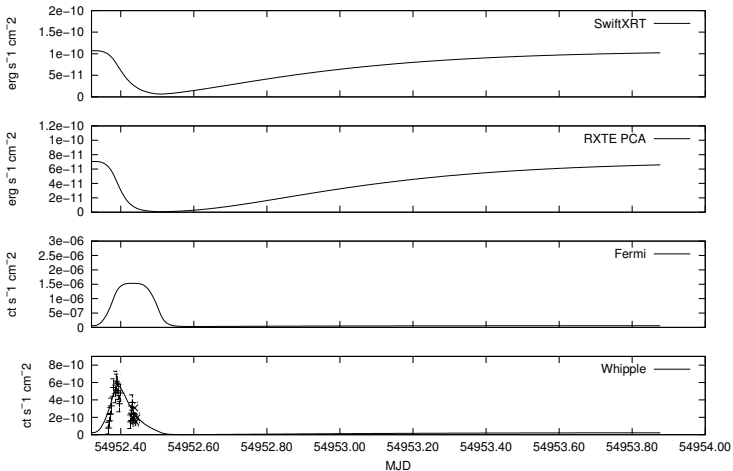


Motivation

- acceleration due to Fermi I happens at finite size shock
 $\Rightarrow R_{acc} \ll R_{rad}$
- acceleration efficiency should depend on distance to shock
- for energies with $t_{cool} < lct$ the blob simply isn't homogenous
 - ▶ example values from Abdo et al. (2011):
 $B = 0.015 \text{ G}, R = 1.3 \cdot 10^{17} \text{ cm}, \gamma_{max} = 1.5 \cdot 10^7$
 - ▶ $t_{cool} = 5.7 \cdot 10^4 \text{ s} \ll t_{esc} = 4.3 \cdot 10^6 \text{ s}$
- compute *multiple shock*-scenarios
- homogenous models constrain time variability to $\Delta t > R_{rad}/c \Rightarrow$
 inhomogenous models allow shorter timescales while preserving causality



Lightcurves for EC orphan flare.



Lightcurves for EC orphan flare.